

Enhanced Photodetection Performance of an In Situ Core/Shell Perovskite-MoS₂ Phototransistor

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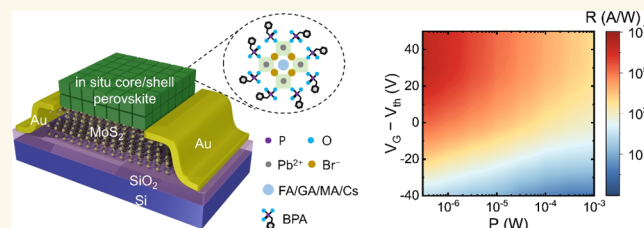
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ABSTRACT: While two-dimensional transition metal dichalcogenides (TMDCs)-based photodetectors offer prospects for high integration density and flexibility, their thinness poses a challenge regarding low light absorption, impacting photodetection sensitivity. Although the integration of TMDCs with metal halide perovskite nanocrystals (PNCs) has been known to be promising for photodetection with a high absorption coefficient of PNCs, the low charge mobility of PNCs delays efficient photocarrier injection into TMDCs. In this study, we integrated MoS₂ with in situ formed core/shell PNCs with short ligands that minimize surface defects and enhance photocarrier injection. The PNCs/MoS₂ heterostructure efficiently separates electrons and holes by establishing type II band alignment and consequently inducing a photogating effect. The synergistic interplay between photoconductive and photogating effects yields a high responsivity of 2.2×10^6 A/W and a specific detectivity of 9.0×10^{11} Jones. Our findings offer a promising pathway for developing low-cost, high-performance phototransistors leveraging the advantages of two-dimensional (2D) materials.

KEYWORDS: heterostructure phototransistor, perovskite nanocrystal, transition metal dichalcogenide, photoconductive effect, photogating effect



Two-dimensional (2D) transition metal dichalcogenides (TMDCs) are promising materials for use in a variety of optoelectronic device applications due to their tunable band gap, simple fabrication, intriguing excitonic properties, and high carrier mobility.^{1–4} Therefore, there have been extensive efforts to exploit these outstanding properties of 2D TMDCs for optoelectronics, especially photodetectors.^{5–10} Photodetectors are basic components of optoelectronics and have applications in various fields such as thermal imaging cameras, optical communications, machine vision technology, and more. However, 2D TMDCs have low light absorption due to their atomic-level thinness, posing a limitation for sensitive photodetection.¹¹

To enhance the light absorption of atomically thin TMDCs, various strategies have been proposed, such as designing plasmonic structures,¹² fabricating in-plane p–n junctions through electrostatic gating,¹³ modulating electronic structures through 2D/2D heterostructures,^{14–17} and integrating highly optically absorbing materials.^{18–21} Among these strategies, heterostructures incorporating lead halide perovskite materials

known for their high optical efficiency, tunable band gap, and solution processability have received considerable attention.^{18–20,22} In particular, recent studies have reported significant performance improvements using perovskite nanocrystals (PNCs), which are zero-dimensional (0D) structures of a perovskite. Light absorption can be maximized through the 0D structure of PNCs, but there is a trade-off of degraded carrier mobility and lower charge injection into TMDC layers, caused by the insulating organic ligands surrounding PNCs.^{11,23,24}

To address the limitations of conventional PNCs and improve the performance of TMDC photodetectors, we integrated MoS₂ phototransistors with in situ formed core/

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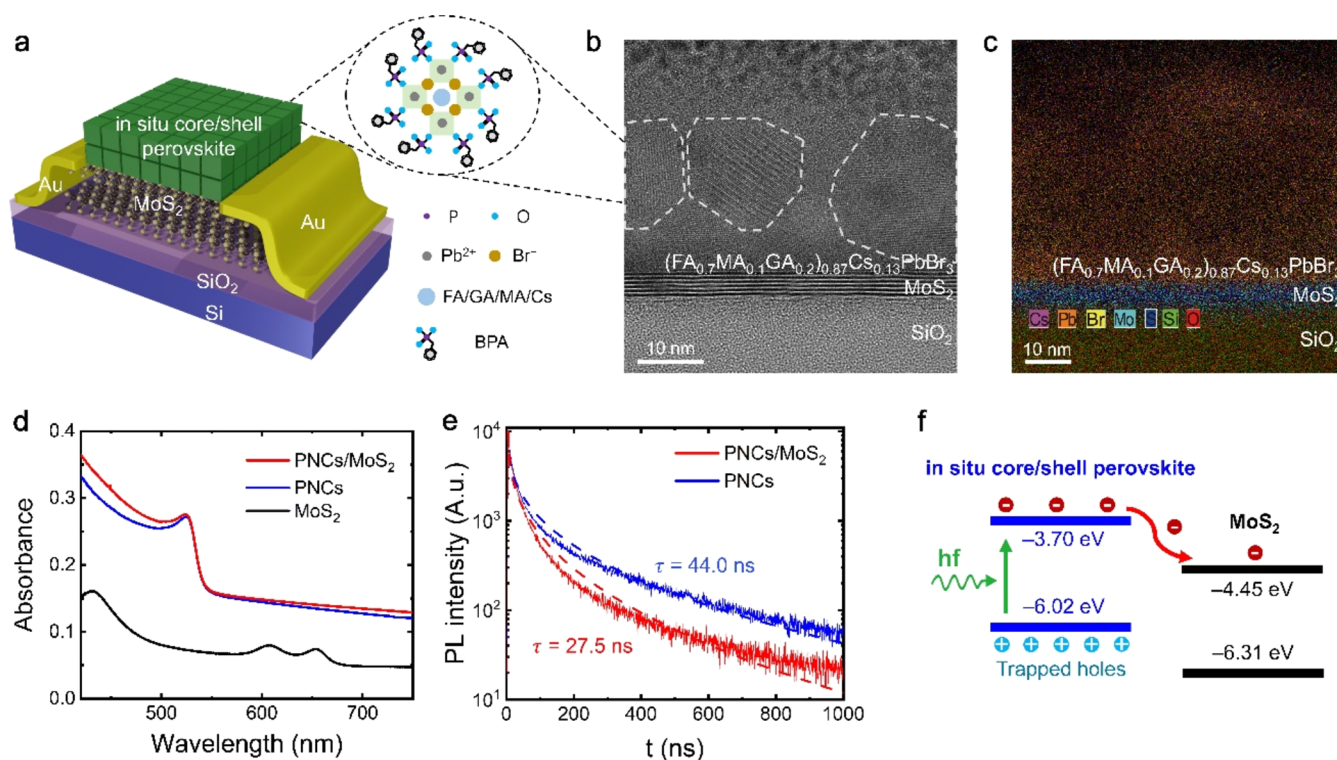


Figure 1. Characteristics of the device. (a) Schematic of PNCs/MoS₂ phototransistor. (b) Cross-sectional TEM image and (c) EDS mapping image of PNCs/MoS₂ phototransistor. (d) UV-vis absorption spectra image of the MoS₂, PNCs, and the PNCs/MoS₂ phototransistor. (e) TRPL image of the MoS₂ phototransistor before and after spin-coating the in situ core/shell perovskite. (f) Band alignment of PNCs/MoS₂ phototransistor.

shell PNCs. As some of us have previously reported,²⁴ in situ core/shell perovskites are created by splitting three-dimensional (3D) perovskite films using benzylphosphonic acid (BPA), which significantly reduces the size of the ligands and minimizes defects. This leads to enhanced carrier mobility of the PNCs and rapid charge injection into TMDC layers.²⁴ Due to these properties, our PNCs/MoS₂ phototransistors exhibited high values of photodetection figures of merit, i.e., a maximum responsivity of 2.2×10^6 A/W and a maximum specific detectivity of 9.0×10^{11} Jones, highlighting its enhanced performance as a phototransistor. We also investigated the photocurrent generation mechanisms of the PNCs/MoS₂ phototransistor, namely, photoconduction and photogating, and discovered their effects on photocurrent by manipulating the gate voltage, ultimately identifying the conditions that yield the best performance. Our research contributes to a deeper understanding of the photocurrent generation phenomena in TMDC phototransistors and offers a promising approach to the study of high-performance phototransistors based on 2D materials.

RESULTS AND DISCUSSION

The device schematic of the PNCs/MoS₂ phototransistor is shown in Figure 1a. In situ formed core/shell PNCs are integrated as a photosensitive layer on top of the MoS₂ field effect transistors (FETs). Details about the device fabrication can be found in the Supporting Information. We examined the PNCs/MoS₂ heterostructure through a transmission electron microscope (TEM) to verify that the PNCs layer is well synthesized on top of the MoS₂ layer. The cross-sectional TEM image of Figure 1b shows the core/shell structure of the PNCs with a cubic crystalline core with a lattice distance of

~ 5.8 Å, identical to our previously reported PNCs synthesized on other substrates.²⁴ A clean interface between the PNCs and the MoS₂ layer was confirmed, which is crucial for the photodetection performance of the device. In addition, the energy-dispersive X-ray (EDS) mapping image of Figure 1c validates the chemical compositions of the PNCs and MoS₂ in the heterostructure. Note that the typical exciton diffusion length in PNCs exceeds several hundred nanometers,²⁵ and the thickness of the PNCs layer (~ 60 nm, Figure S2) is suitable for excitons generated in the uppermost layer of the PNC to migrate to the PNCs/MoS₂ interface before recombination within the PNCs layer (Figure S3).

The role of PNCs as photosensitive layers is to efficiently absorb light and inject photogenerated carriers into the MoS₂ channel. We performed optical characterizations to confirm that the PNCs layer is effective as a photosensitive layer, as shown in Figure 1d,e. First, the ultraviolet-visible (UV-vis) spectra in Figure 1d demonstrate that the light absorption in the PNCs/MoS₂ heterostructure is significantly increased compared to that of MoS₂ alone. The enhancement of light absorption can be attributed to the high light absorption coefficient of lead halide perovskites.²² Moreover, the absorbance spectra of the PNCs/MoS₂ heterostructure and isolated PNCs show similar intensities, indicating that the light absorption of the MoS₂ layer is negligible in the PNCs/MoS₂ heterostructure. It shows that lights will mainly be absorbed by the PNCs photosensitive layer in the PNCs/MoS₂ phototransistor. This can also be seen by comparing the absorbance peak after the integration of the PNCs layer on MoS₂. The absorption spectrum of the MoS₂ shows two excitonic peaks between 600 and 700 nm which originate from the A and B direct excitonic transitions, and another excitonic peak near

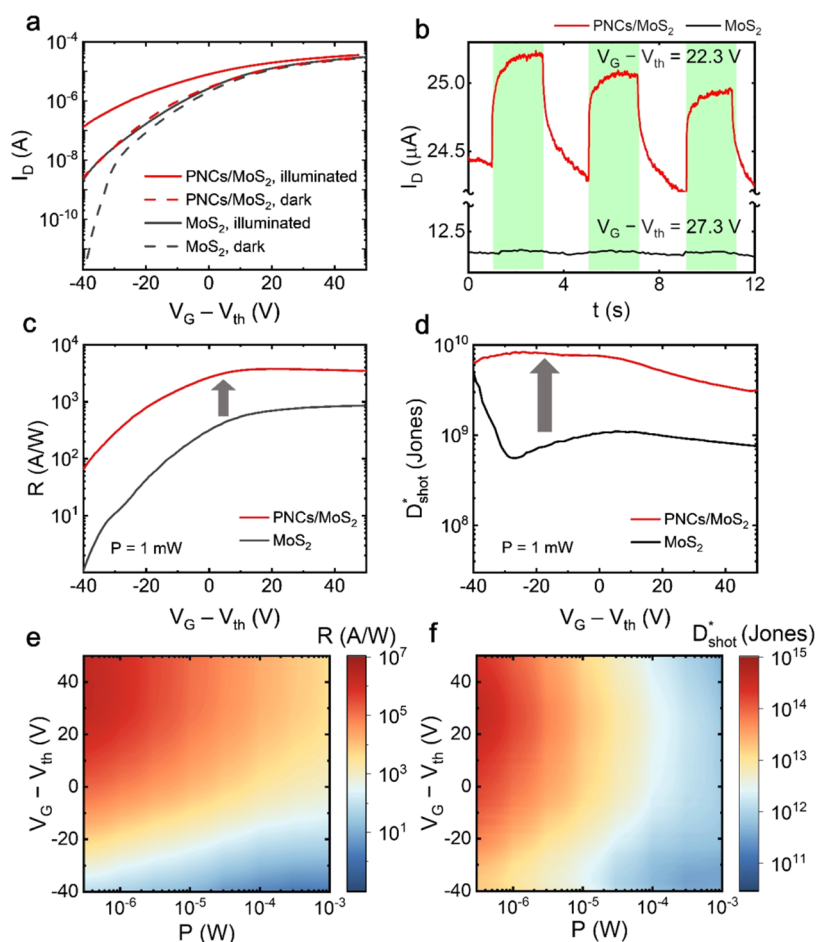


Figure 2. Optoelectrical properties of the device. (a) Transfer curves and (b) photoresponses of PNCs/MoS₂ and MoS₂ phototransistors at a laser power of 1 mW. (c) Responsivity and (d) specific detectivity derived from shot noise of PNCs/MoS₂ and MoS₂ phototransistors at a laser power of 1 mW. (e) Responsivity and (f) specific detectivity derived from shot noise of the PNCs/MoS₂ phototransistor in different gate voltage and laser power. All measurements were conducted at $V_D = 0.1$ V.

400 nm which is related to C-exciton.^{26,27} On the other hand, the absorption spectrum of the PNCs/MoS₂ heterostructure lacks the excitonic peaks characteristic of MoS₂. Instead, an absorption peak at 540 nm appears, corresponding to the distinct band edge absorption peak of PNCs, as confirmed in our previous report.²⁴ This observation further confirms the dominant role of the PNCs layer in light absorption.

Time-resolved photoluminescence (TRPL) analysis was performed to verify that photogenerated electrons are transferred to the MoS₂ channel from the photosensitive layer (the PNCs) in Figure 1e. The PL intensity decay was fitted to the stretched exponential function $I(t) = I(0) \exp[-(t/\tau^*)^\beta]$. In this equation, β is the stretching parameter, and τ^* is the time constant. The stretched exponential function is commonly used to model the time relaxation behaviors in disordered systems, such as the TRPL decay.^{28,29} The PL lifetime (τ value) can be obtained by the equation $\tau = \int_0^\infty \exp[-(t/\tau^*)^\beta] dt = \frac{\tau^*}{\beta} \Gamma\left(\frac{1}{\beta}\right)$, where Γ is the gamma function. The PL lifetime observed for the PNCs/MoS₂ heterostructure (~ 27.5 ns) is shorter compared to that of the PNCs (~ 44.0 ns), suggesting that radiative recombination decreased within the PNCs/MoS₂ heterostructure. This is attributed to the efficient charge transfer of photogenerated carriers from the PNCs to the MoS₂ layer.

Photogenerated carrier transfer from the PNCs to the MoS₂ can be explained by the type II band alignment of the PNCs/MoS₂ heterostructure as shown in Figure 1f. The band alignments are confirmed by Kelvin probe force microscopy and ultraviolet photoelectron spectroscopy (Figure S6).³⁰ When electron–hole pairs are generated within the PNCs, electrons from the PNCs tend to transfer to the MoS₂ due to the lower conduction band minimum level of the MoS₂ compared with that of PNCs. On the other hand, holes remain trapped within the PNCs, since the valence band of PNCs has a higher energy level than that of the MoS₂. Here, the rate of electron transfer (κ) can be calculated by the equation

$$\kappa = \frac{1}{\tau_{\text{PNCs/MoS}_2}} - \frac{1}{\tau_{\text{PNCs}}} \quad (1)$$

where $\tau_{\text{PNCs/MoS}_2}$ and τ_{PNCs} are the PL lifetime of PNCs/MoS₂ heterostructure and PNCs layer, respectively.²³ In our study, κ was found to be $\sim 10^7$ s⁻¹, indicating the rapid transfer of photogenerated electrons from the PNCs to the MoS₂, compared to previously reported devices.^{23,31} These charge transfer behaviors of the photogenerated electrons and holes result in both a photoconductive effect and a photogating effect. The photoconductive effect stems from the increased carrier density of the n-type MoS₂ channel due to the

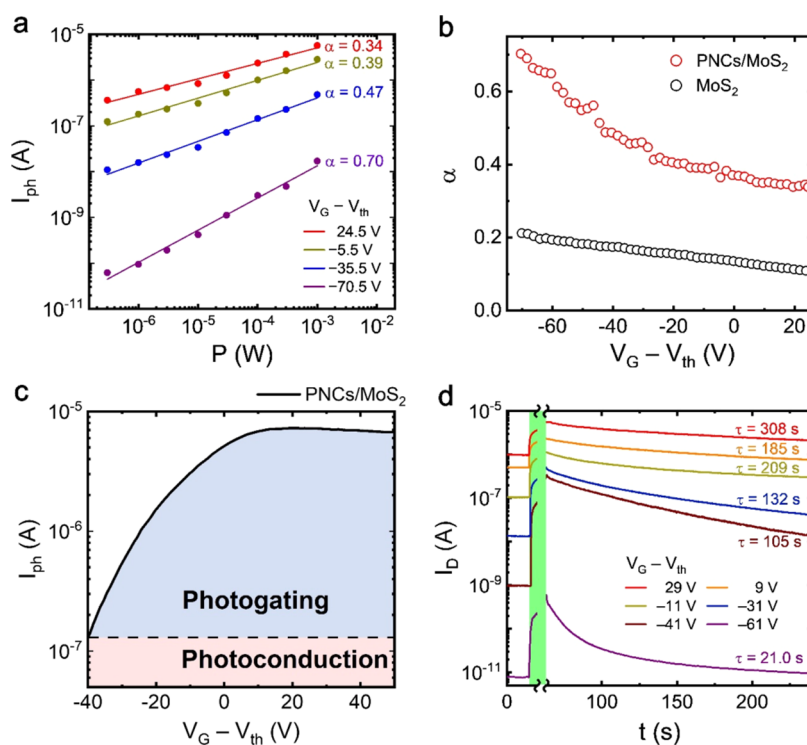


Figure 3. Photocurrent generation mechanisms of the PNCs/MoS₂ phototransistor. (a) The photocurrent curves fitted to a simple power law for laser power. (b) The gate voltage dependence of α . (c) Gate dependence of the photocurrent. (d) Response time of the device as a function of gate voltage.

transferred electrons. On the other hand, the photogating effect emerges through the accumulation of holes in the PNCs, leading to a positive gating effect that subsequently amplifies the photocurrent within the MoS₂ channel.⁵

To confirm the improvement of the photodetection performance due to the integration of the PNCs, we compared the optoelectronic properties of the PNCs/MoS₂ phototransistor with those of the MoS₂ phototransistor without the PNCs layer using a 532 nm laser with an optical power of 1 mW. Note that a moderate optical power of 1 mW was selected to maintain the stability of the phototransistor, preventing any potential damage due to excessively strong light illumination. The transfer characteristics, i.e., drain current (I_D) versus overdrive voltage ($V_G - V_{th}$, V_{th} is threshold voltage) of the devices are shown in Figure 2a. In this figure, the drain currents were plotted against the overdrive voltage to compare the currents at the same gate voltage above the threshold voltage. Here, the photocurrent (I_{ph}) is defined by the following equation $I_{ph} = I_{illum} - I_{dark}$, where I_{illum} and I_{dark} are each current under illuminated and dark conditions. As shown in Figure 2a, the PNCs/MoS₂ phototransistor exhibits higher I_{ph} and I_{illum} values compared to those of the MoS₂ phototransistor. This amplified photoresponse arises from the increased light absorption in the PNCs layer coupled with the synergistic interaction of the photoconductive and photogating effects due to the type II band alignment of the PNCs/MoS₂ heterostructure (Figure 1f). It should be noted that the increase in dark current at low gate voltage, reaching near a nanoampere, is attributed to electron transfer due to the energy level difference between the PNCs and MoS₂, as well as a leakage current of 0.2 nA through the PNCs layer (Figure S8). However, in most gate regions, the leakage current through the PNCs is negligible compared

to the photocurrent in the range of several μ A, and it does not affect the operation mechanism and the performance of the device.

Enhanced photoresponse of the PNCs/MoS₂ phototransistor is more distinctly evident in the dynamic photoswitching behavior in Figure 2b. A light pulse with a duration of 2 s was applied every 4 s to compare the photoswitching behavior at the similar value of $V_G - V_{th}$. As a result, the PNCs/MoS₂ phototransistor showed clear photoswitching behavior, which was superior to the case of MoS₂.

Figure 2c,d shows the responsivity and specific detectivity of the PNCs/MoS₂ and the MoS₂ phototransistors. Responsivity (R) can be obtained by the following equation $R = I_{ph}/P_{eff}$, where P_{eff} is the effective incident power which can be calculated by multiplying the optical power of the laser by the ratio of the MoS₂ channel area to the laser beam spot size. Also, based on the shot noise approximation, specific detectivity (D^*) can be obtained by the following equation $D^*_{shot} = R\sqrt{A}/\sqrt{2eI_{dark}}$, where e is the elementary charge and A is the area of a photosensitive region. As shown in Figure 2c,d, the PNCs/MoS₂ phototransistor demonstrated higher responsivity and specific detectivity compared with those of the MoS₂ phototransistor for all V_G ranges.

In Figure 2c, the responsivity of both PNCs/MoS₂ and the MoS₂ phototransistors increases at higher V_G and saturates near $V_G - V_{th}$ of ~ 10 V. This trend is frequently observed in photodetectors that are primarily driven by the photogating effect.³² The photogating effect of the PNCs/MoS₂ heterostructure is attributed to the trapped holes in the PNCs and the photogating effect effectively increases the photocurrent, enabling high responsivity. On the other hand, the specific detectivity of the PNCs/MoS₂ phototransistor tends to decrease at higher gate voltages as shown in Figure 2d. This

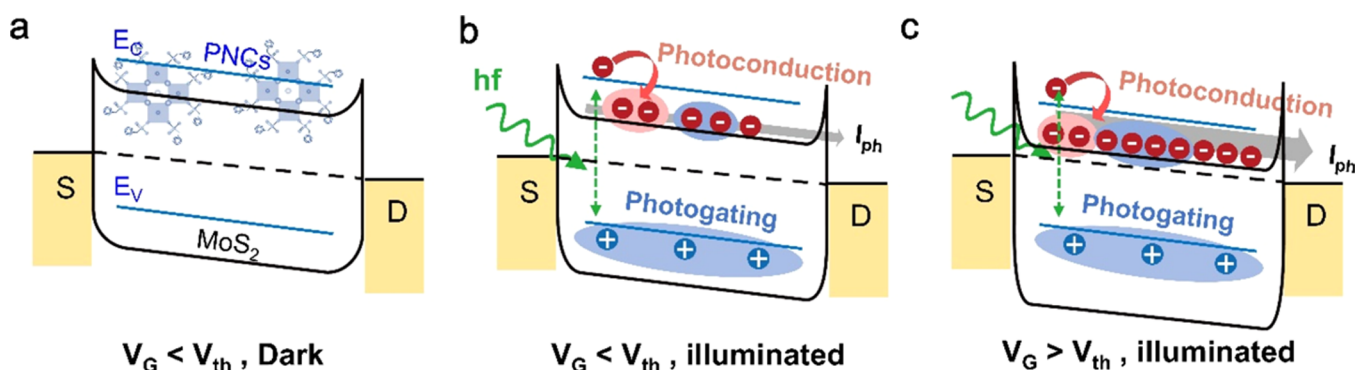


Figure 4. Band diagram illustrating the operational modes of the PNCs/MoS₂ phototransistor as a function of gate voltage and light conditions: (a) $V_G < V_{th}$, dark, (b) $V_G < V_{th}$, illuminated, and (c) $V_G > V_{th}$, illuminated.

trend is due to detectivity being defined as the responsivity divided by the square root of the dark current, which in turn increases at higher gate voltages. Note that in the case of MoS₂ phototransistors, high detectivity was observed at a very low gate voltage of -40 V, which is attributed to its much lower dark current compared to the case of PNCs/MoS₂ phototransistors.

To further analyze the photodetection performance of PNCs/MoS₂ phototransistors, contour plots depicting the responsivity and specific detectivity as functions of $V_G - V_{th}$ and incident light power are presented in Figure 2e,f. Responsivity increased with increasing V_G , while the specific detectivity tended to peak at intermediate V_G between 10 and 30 V. Furthermore, both responsivity and detectivity showed a tendency to increase as the light intensity decreased. Note that the increase in responsivity under weak light conditions is also a major indication of the photogating effect.³³ The maximum responsivity and specific detectivity within the measurement range for our phototransistors were 2.2×10^6 A/W and 9.0×10^{11} Jones respectively, which are significantly higher than previously reported values for the photodetectors utilizing similar device structures. Note that specific detectivity in Figure 2d,f is calculated based on shot noise. To calculate the accurate value of specific detectivity,³⁴ low-frequency noise measurements were performed (Figure S15). The comparison to the previous report is discussed in more detail in Figure 5.

The photogeneration mechanisms in photodetectors made of 2D-layered materials primarily involve photoconduction and photogating effects.³³ Photoconduction is a phenomenon in which carrier density increases by photogenerated carriers, thereby directly contributing to the increase in channel conductance. As the number of photons and carriers is proportional to the incident power, photocurrent has a linear dependence on incident power ($I_{ph} \propto P^\alpha$, $\alpha = 1$). On the other hand, photogating occurs when carriers generated by light become trapped at interface traps, acting as a gating source to the channel. Charge trapping on the SiO₂ surface due to water molecules³² and spatial potential fluctuations caused by local strains between SiO₂ and 2D materials³⁵ are such examples of interface traps. For our device, MoS₂ and PNCs form a type II band alignment, so accumulated holes in the PNCs are the main cause of the photogating effect. In the case of photogating, the photocurrent follows a fractional power law dependence of incident power ($I_{ph} \propto P^\alpha$, $0 < \alpha < 1$).^{36,37} If α largely varies from 1, the nonlinearity of photocurrent to the incident power is large, which indicates that photogating acts as a dominant mechanism for photocurrent generation.^{11,23}

The nonlinearity of the photocurrent with respect to incident light power was calculated to observe the photogating effect at different gate voltages, as shown in Figure 3a,b. The trend lines in Figure 3a display the power law fitting $I_{ph} \propto P^\alpha$ for different gate voltages with the slope of the trend lines representing α . As the gate voltage increased, it was observed that α decreased from 0.70 to 0.34 and that the nonlinearity of the photocurrent dependence on incident power increased. It can be seen that α decreased as the gate voltage increased in Figure 3b. This indicates that the influence of the photogating effect contributing to the photocurrent increases as the gate voltage increases. This trend is visually confirmed in Figure 3c. Considering that photoconduction applies a constant DC offset to the photocurrent, the remaining photocurrent can be attributed to photogating. In other words, the gate dependence of the photocurrent allows us to confirm the gate dependence of the photogating effect, which is consistent with the trend observed in Figure 3b. It should be noted that the photogating effect can be interpreted as a shift of the threshold voltage, resulting in a similarity between the gate dependence of the photocurrent (Figure 3c) and the transconductance (Figure S11). Also, for reference, the value of α maintained a very low value below 0.25 across all V_G regions for MoS₂, indicating that photogating had a significant influence on MoS₂. This well agrees with the reported persistent photoconductivity phenomena of MoS₂ arising from various factors such as defects, surface adsorbates, and MoS₂/SiO₂ trap states.³⁸

Another distinctive feature between photogating and photoconductive effects is the response times. Photoconduction stops generating electron–hole pairs when the light is turned off, leading to an almost immediate return to the dark current and thus to a short response time. In contrast, the photogating effect, even after the light is turned off, requires a large amount of time for the trapped charges to be detrapped before recovering to the dark current, resulting in a much longer response time.³³ Figure 3d shows the response time of the PNCs/MoS₂ phototransistor to observe the photogating effect at different gate voltages. Fall time is defined as the time it takes for I_{illum} to return to $0.1 \times I_{illum} + 0.9 \times I_{dark}$. It can be observed that the fall time increases as the gate voltage becomes higher. This also supports the fact that the influence of photogating to contribute to the photocurrent increases as the gate voltage increases. Also, it is worth mentioning that as gate voltage increases, the photocurrent increases and fall time also increases, so the optimal gate voltage condition should be determined for the specific purpose.

The observation can be understood through the energy band diagrams in Figure 4. Figure 4a depicts the energy band diagram of the phototransistor in the dark, with a gate voltage less than the threshold voltage ($V_G < V_{th}$). In this state, the carrier concentration within the MoS₂ channel is low. Figure 4b explains the situation when the device is illuminated with light at a gate voltage of less than the threshold voltage. The PNCs absorb the photons and generate electron–hole pairs. Electrons transfer to the MoS₂ layer and lead to an increased number of charge carriers, which is the photoconductive effect, marked in red. Concurrently, holes are confined within the PNCs, marked in blue, creating a photogating effect. Both effects synergistically contribute to the increase in the photocurrent. In Figure 4c, the condition changes such that the device is illuminated while the gate voltage is greater than the threshold voltage ($V_G > V_{th}$). Here, the channel already has an abundance of charge carriers due to the higher gate voltage. While the photogating effect is modulated by the gate voltage, the photoconductive effect remains relatively gate-independent.³⁶ Therefore, the photogating effect, indicated in blue, becomes more pronounced, contributing a larger proportion to the total photocurrent. Figure S12 provides a more detailed view of the gate dependence of the photogeneration mechanisms.

Figure 5 summarizes the responsivity and specific detectivity of various photodetectors based on 2D materials reported to

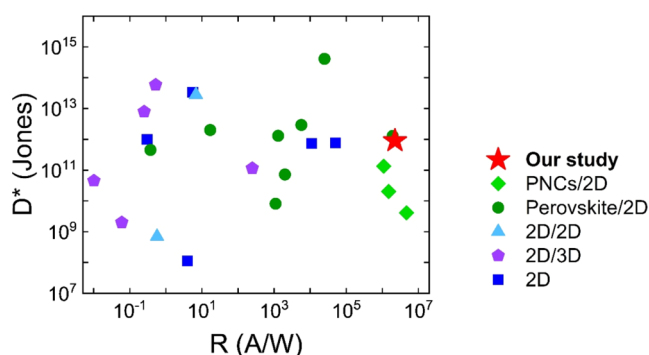


Figure 5. Photodetection performance comparison of various photodetectors based on 2D materials,^{39–43} 2D/2D heterostructures,^{44,45} 2D/3D heterostructures,^{46–50} perovskite film/2D heterostructures,^{18,19,51–56} and perovskite nanocrystals/2D heterostructures.²³

date (Table S1 and Figure S16). Due to the extremely low dark current, 2D materials alone can exhibit high photodetection performance. However, various heterostructures have been explored to enhance its low light absorption. The 2D/2D heterostructure and perovskite film/2D heterostructure are prominent examples, but they have encountered their respective limitations and have not shown significant performance improvement compared to 2D materials. Furthermore, the PNCs/2D heterostructure utilizing 0D perovskite effectively enhanced responsivity but resulted in decreased specific detectivity due to high dark current. Unlike these heterostructures based on low-dimensional perovskite materials, in situ core/shell perovskite achieved significantly improved responsivity followed by improved specific detectivity. Since the in situ core/shell perovskite is synthesized from a thickness-controllable 3D perovskite thin film followed by an in situ splitting process, significantly increased absorbance in the active layer could be achieved by utilizing a thickness far

greater than that of 2D perovskites or colloiddally synthesized PNCs. Furthermore, the fast charge transport from the absence of long alkyl ligands and passivated surface trap states in core/shell structure collectively minimize carrier loss due to trap-mediated nonradiative recombination, a common issue in the extended charge transport and injection pathways toward the TMDC layer.

In future research, efforts to enhance the commercial potential of the device could be considered. First, employing MoS₂ produced via large-area film deposition, rather than mechanically exfoliated MoS₂ flakes, may facilitate broader production capabilities. Second, coating PNCs through thermal evaporation could simplify mass production. Lastly, encapsulating the PNCs could significantly improve their stability.

CONCLUSIONS

In conclusion, we have successfully developed in situ core/shell perovskite nanocrystals (PNCs)/MoS₂ heterostructure phototransistors with the aim of enhancing the photodetection performance. Through optical characterizations, we observed efficient light absorption in the PNCs layer along with efficient charge transfer from the PNCs to the MoS₂. The PNCs/MoS₂ phototransistors showed increased photocurrent compared to the MoS₂ phototransistors without PNCs. Moreover, the photodetector's responsivity and specific detectivity were also enhanced, reaching maximum values of 2.2×10^6 A/W and 9.0×10^{11} Jones, respectively. To elucidate the photocurrent generation mechanism, we systematically varied the gate voltage and the incident light power to the PNCs/MoS₂ phototransistors and observed the relationship between photocurrent and incident light power to identify two distinct phenomena of photogating and photoconductive effects. These results indicate that while photogating and photoconductive effects synergistically contribute to generating photocurrent, photogating becomes more dominant at higher gate voltages. We believe that our perovskite-MoS₂ heterostructure phototransistors demonstrate the potential of 2D material-based photodetectors, especially when they are integrated with an efficient photosensitive layer.

METHODS

MoS₂ FET Fabrication. MoS₂ flakes were mechanically exfoliated from bulk crystals and transferred to a 270 nm SiO₂/p++ Si substrate, with flake selection and thickness measurement conducted using an optical microscope and an AFM system (NX-10, Park Systems). After spin-coating of methyl methacrylate and poly(methyl methacrylate) resist layers on the MoS₂ flakes, source and drain electrodes were patterned using an electron-beam lithography system (JSM-6510, JEOL). Then, a 50 nm Au layer as electrodes was deposited using an electron-beam evaporator (KVE2004 L, Korea Vacuum Tech). The completed MoS₂ FETs were annealed at 150 °C for 1 h to enhance their electrical properties before measurement.

Integration of In Situ Core/Shell Perovskite of MoS₂ FET. To prepare the precursor solution of the mixed cation ((FA_{0.7}MA_{0.1}GA_{0.2})_{0.87}Cs_{0.13}PbBr₃), specific stoichiometric amounts of FABr, MABr, GABr, CsBr, and PbBr₂ were dissolved in dimethyl sulfoxide (DMSO) at a concentration of 0.6 M. Additionally, BPA was added at a concentration of 10 mol % relatives to PbBr₂. This solution was then stirred continuously overnight in a N₂-filled glovebox, ensuring thorough mixing and dissolution before its subsequent use. To synthesize in situ core/shell perovskite films, the substrate was first subjected to ultraviolet-ozone treatment to increase its hydrophilicity. The substrates were then moved to a N₂-filled glovebox. Here, metal halide perovskite (MHP) films, with a thickness of about 60 nm, were

formed using a spin-coating technique at 6000 rpm. This process involved the A-NCP method, where during the second stage of spinning, a solution of 2,2',2''-benzene-1,3,5-triyltris(1-phenyl-1H-benzimidazole) (TPBi) in chlorobenzene (CB) was applied to the perovskite film. Following this, a solution of BPA in tetrahydrofuran (THF) was spread over the perovskite layer. The assembly was then allowed a reaction time of 15 s before being spun dry immediately afterward.

Optical and Electrical Characterization. The PL and Raman spectra were obtained through a confocal imaging system (Xper Raman 200, Nanobase) with an incident laser beam of wavelength 532 nm. The TRPL spectra were obtained through a confocal fluorescence lifetime image microscope system (FlouTime 300, PicoQuant), and the absorption spectra were obtained by using a UV–visible spectrophotometer (V-730, JASCO). The KPFM method is conducted using the AFM system (NX-10, Park Systems), and the UPS method is conducted using an X-ray photoelectron spectrometer (Versaprobe III, Ulvac-PHI). The TEM images were obtained by high-resolution transmission electron microscopy (FEI Titan 80–30 TEM, Thermo Fisher). Monolayer MoS₂ synthesized via chemical vapor deposition (CVD) was used for TRPL and UV–visible spectroscopy. The electrical characteristics of FETs were measured by using a probe station (M6VC, MSTECH) and a semiconductor parameter analyzer (Keithley 4200). The photoresponses of FETs were measured under laser (MDE5240 V) illumination of wavelength 520 nm. The laser beam was globally illuminated to phototransistors with a few millimeters in diameter. All of the characterizations were performed at room temperature. To minimize unwanted ambient air exposure, PNCs/MoS₂ phototransistors were moved from an N₂-filled glovebox to characterization equipment using metal containers. These containers are designed to keep a vacuum inside and shield the devices from external degrading factors such as humidity, various gases in the environment, and light exposure.

ASSOCIATED CONTENT

Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acsnano.4c02775>.

Additional details on device fabrication process, thickness and optical characterization of in situ core/shell perovskite, MoS₂, PL lifetime of PNCs/MoS₂ heterostructure with different thickness, KPFM and UPS characterization of PNCs/MoS₂ heterostructure for calculating the energy band alignments, electrical characterization of PNCs and MoS₂ FET, dynamic photoswitching behavior of MoS₂ FET, detailed explanation of photogating and photoconductive effect on PNCs/MoS₂ phototransistors, stability and noise of PNCs/MoS₂ phototransistors, and comparison of photo-detection figure of merits with other reports (PDF)

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Notes

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